A holomorphic approach to fivebranes

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Most work is joint with PhD student Surya Raghavendran at Perimeter Institute.



And some of it with Ingmar Saberi at LMU.



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The vertex algebra popping out of the AGT correspondence has the interpretation as the algebra of operators on the Ω -deformed fivebrane theory. In the bulk of the talk we will follow a similar approach before doing the equivariant localization.

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The worldvolume theory associated to fivebranes is a six-dimensional $\mathcal{N} = (2,0)$ supersymmetric theory.

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theory becomes an effective QM system. The state space becomes the $U(1) \times U(1)$ -equivariant cohomology of the moduli space of instantons.

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Fivebranes (which are six-dimensional objects before localization) become two-dimensional objects wrapping

$$0 \times C \subset \mathbf{R} \times \mathbf{T}^* C$$

which are chiral defects in the five-dimensional theory.

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- We don't perform equivariant localization (Ω -background) with respect to the rotation action by $U(1) \times U(1)$.

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There are basically two classes of twists characterized by the number of invariant directions:

- The minimal twist, exists on any complex three-fold Z^* .
- The non-minimal twist, this exists on any manifold of the form $C \times M^4$.

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Degree zero $\overline{\partial}$ -cohomology is

$$H^0(\mathbf{C}^3 - 0, \mathcal{O}) = \mathcal{O}(\mathbf{C}^3).$$

The 'singular' part of the OPE lives in $H^2(\mathbb{C}^3 - 0, \mathcal{O})$.

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Theorem (Saberi-W.)

The holomorphic twist of the abelian six-dimensional $\mathcal{N} = (2,0)$ theory has space of fields consisting of:

- a three-form $\alpha \in \Omega^{2,1}(Z)$ which satisfies the constraint $\partial \alpha = 0$.
- a symplectic-valued section $\phi \in \Omega^{0,1}(Z, K^{1/2}) \otimes \mathbb{C}^2$.

The 'action functional' is

$$\int_{Z} \alpha \overline{\partial} \partial^{-1} \alpha + \int_{Z} \phi \overline{\partial} \phi.$$

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The moduli space of solutions, modulo gauge symmetries, is a bundle over the intermediate Jacobian variety of Z.

We proceed similarly to Costello.

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The idea is that the Hilbert space of the five-dimensional theory on X in the holomorphic twist is the space of de Rham forms on $Bun_G(X)$, which further deforms to de Rham cohomology in the topological twist. One form of evidence we have is that under the equivariant localization the algebra $\mathcal{A}_{\mathfrak{gl}(1)} = \mathcal{A}_{\mathfrak{gl}(1)}(\mathbf{C}^2)$ becomes a familiar one.

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We now turn to our holographic approach to non-abelian fivebranes.

In some limit, pieces of string / M theory are supposed to be described by supergravity.

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Roughly, fields are the cyclic cohomology of category of coherent sheaves on a CY—via HKR we can identify this with polyvector fields and we recover Beltrami equation

$$\overline{\partial}\mu + \frac{1}{2}[\mu,\mu] = 0,$$

as part of the equations of motion.

With Raghavendran and Saberi we described twists of 11d supergravity.

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- $\mu \in \Pi\Omega^{\bullet}(\mathbf{R}) \otimes \mathrm{PV}^{1,\bullet}(X)$ required to be divergence-free and locally constant along **R**.
- $\gamma \in \Omega^{\bullet}(\mathbf{R}) \otimes \Omega^{1,\bullet}(X).$

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 \star <u>Caveat</u>: twisted theory carries only an overall **Z**/2 grading.

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Theorem (Raghavendran-Saberi-W.) The algebra of local operators on $\mathbf{R} \times \mathbf{C}^5$ is equivalent to

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We now expand on the program of twisted holography.

Ordinary Koszul duality for associative algebras makes its appearance in QFT when studying (topological) line operators.

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$\operatorname{Obs}\otimes B.$

Couplings are determined by MC elements $\alpha \in Obs \otimes B$.

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Upshot: Obs[!] is the algebra describing the *universal* line defect.

Example:

Example: Line operators in CS theory on $\mathbf{R} \times M = \mathbf{R} \times \mathbf{R}^2$.

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To couple line operator with algebra of operators ${\cal B}$ we must prescribe an algebra map

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 $\rightsquigarrow B$ is a \mathfrak{g} -module!

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Only understand in special examples!

Steps to characterize 'universal defect':
• Identify observables Obs_{sugra} of the bulk supergravity theory as a factorization algebra on M_{bulk} .

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- **2** Pick a submanifold $Z \subset M_{bulk}$ where defect will be supported ('wrap'). Then 'restrict' $Obs_{sugra}|_Z$ to a factorization algebra on Z.
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Back to fivebranes,

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Back to fivebranes, we look at twisted supergravity on the eleven-manifold

$$M_{bulk} = \mathbf{R} \times \mathrm{Tot}(V \to Z)$$

where: Z is a complex three-fold and $V \rightarrow Z$ is a rank two holomorphic vector bundle satisfying $\wedge^2 V \cong K_Z$. The location of the branes is along the zero section

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Expectation: the universal theory along the brane = 'theory on stack of $\overline{N} = \infty$ many (twisted) fivebranes' is

$$(\mathrm{Obs}_{sugra}|_Z)^! = (\pi_*\mathrm{Obs}_{sugra})^!$$

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Question: What does twisted holography tell us about

 $Obs_{\mathfrak{gl}(N)}$ or $Obs_{\mathfrak{sl}(N)}$?

We make use of a filtration found by Kac-Rudakov.

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Every V_j is an irreducible E(3|6) representation. Admits elegant description as quotients of certain E(3|6) Verma modules.

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First nontrivial associated graded is the factorization algebra

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Proposition (Raghavendran-W.)

There is a quasi-isomorphism of factorization algebras on the threefold Z (even at the quantum level)

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A good start! What about the next layer?

The next piece has the explicit form

$$\mathcal{F}^{(0)}/\mathcal{F}^{(1)} \simeq \mathcal{C}_{\bullet}(\mathcal{E}(3|6)_c)$$

where $\mathcal{E}(3|6)$ is a sheaf of super dg Lie algebras on threefold Z which enhances Kac's exceptional super Lie algebra E(3|6).

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Conjecture (Raghavendran-Saberi-W.)

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We also have a putative quantization computed by Witten diagrams in twisted supergravity—appears as turning on a cocycle for the Lie algebra $\mathcal{E}(3|6)$.

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Recall that we have understood the equivariant localization as an explicit deformation by a particular superconformal element. Some evidence for this conjecture connects to an original instance of the AGT correspondence.

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Theorem (Saberi-W.)

$$\left(\mathcal{F}^{(0)}/\mathcal{F}^{(1)}, S\right)$$

is equivalent to the vertex algebra of observables of the chiral sector of the Liouville CFT—that is, the Virasoro algebra of a particular central charge.

The Lie algebra $E(3|6) = \mathcal{E}(3|6)|_{\mathbb{C}^3}$ plays a dual role as an infinite-dimensional enhancement of the twisted superconformal algebra.

The global symmetry algebra we consider is

 $\mathfrak{gl}(1) \times \mathfrak{sl}(3) \times \mathfrak{sl}(2) \subset \mathfrak{osp}(6|2),$

with fugacities (q, t_1, t_2, r) .

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For $\mathcal{F}^{(0)}/\mathcal{F}^{(1)}$ we find the closed form character of operators

$$\operatorname{PExp}\left[\frac{q^4(t_1^{-1}+t_1t_2^{-1}+t_2)+q^3(r^2+r^{-2}+1)-q^{7/2}(r+r^{-1})(t_1^{-1}+t_1t_2^{-1}+t_2)}{(1-t_1^{-1}q)(1-t_1t_2^{-1}q)(1-t_2q)}\right]$$

which we expect is the refined supersymmetric index of the 6d $\mathcal{N} = (2,0)$ theory for group SU(2).

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Notice the specialization $q = r^2, t_1 = 1$ is $\operatorname{PExp} \frac{q^2}{1-q} = \operatorname{character}$ of Virasoro vacuum.
Given this decomposition of $\widehat{E(5|10)} = \bigoplus_{j \ge -1} V_j$ into E(3|6)-modules we have a systematic way to compute characters in the holomorphic twist at every level of the filtration.

Specialization $q = r^2, t_1 = 1$ of local character of $\mathcal{F}^{(0)}/\mathcal{F}^{(N-1)}$ agrees with type A_{N-1} \mathcal{W} -algebra vacuum character for every N(and in the large N limit as above).

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Future work is to use holography to describe the quantization of the factorization algebras $\mathcal{F}^{(0)}/\mathcal{F}^{(N-1)}$. Also, can we use this program to give a holomorphically twisted construction of class \mathcal{S} theories?