

Giant Gravitons in Twisted Holography

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Twisting Supersymmetric QFTs

Twisting SQFT: restricting to a cohomology of a supercharge

Physics motivation: produces a protected subsector of SQFT that is easier to study

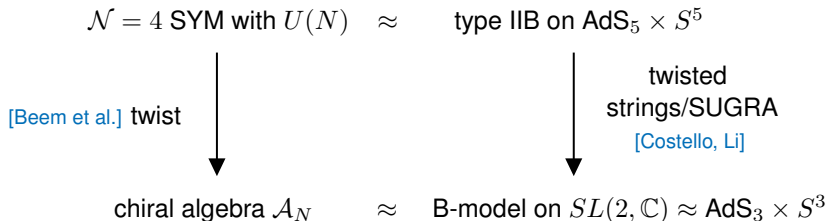
Examples:

- A and B topological twists of 2d $\mathcal{N} = 2$ sigma models [Witten '91]
- holomorphic twist of 4d $\mathcal{N} = 1$ [Johansen '94] [Nekrasov '96] [Costello '13] [Gwilliam, Williams '18] [Saber, Williams '19] [Davide Gaiotto's talk]
- 2d chiral algebra subsector of 4d $\mathcal{N} = 2$ SCFT [Beem, Lemos, Liendo, Peelaers, Rastelli, Rees '13]

Twisted holography: holographic duals of these twists

Twisted Holography

Example: protected subsector of $\text{AdS}_5/\text{CFT}_4$ [Costello, Gaiotto '18]:



Motivation:

- Simplifications
- Connections to math

Twisted Holography

Examples of **twisted supergravity/superstrings**:

- Twisted type IIA and IIB supergravity [Costello, Li '16] [Saberi, Williams '21]
- Twisted 11d supergravity [Saberi, Williams '21] [Eager, Hahner '21]
[Raghavendran, Saberi, Williams '21]
- Twisted S-duality [Raghavendran, Yoo '19]

Other examples of **twisted holography**:

- D2-D4 system in 6d topological strings [Ishtiaque, Moosavian, Zhou '18]
- Subsector of duality between SUGRA on $\text{AdS}_3 \times S^3 \times T^4$ and symmetric orbifold $\text{Sym}^N(T^4)$ [Costello, Paquette '20]
- M2 and M5 branes in an Ω -background [Costello '16 '17]

In this talk

- Review the duality [Costello, Gaiotto '18]

2d chiral algebra $\mathcal{A}_N \approx$ topological B-model on $SL(2, \mathbb{C})$

- Match saddles of determinant correlation functions with D1-brane configurations
 - Spectral curve construction
- Extend the match to non-conformal vacua of the chiral algebra
- Future directions: "bootstrapping" to $AdS_5 \times S^5$, holomorphic twist

Twisted Holography

Twisted holography example [Costello, Gaiotto '18]:

B-model topological strings
on deformed conifold $SL(2, \mathbb{C})$
with coupling N^{-1} \approx large N expansion
of chiral algebra $\mathcal{A}_N =$
gauged $\beta\gamma$ system in adj. of $U(N)$

Simplifications:

- (almost) free field theory computations in the chiral algebra \mathcal{A}_N
- D1-branes are **holomorphic curves** in $SL(2, \mathbb{C})$
- dependence on t'Hooft coupling drops out

Chiral algebra \mathcal{A}_N

The chiral algebra subsector of $\mathcal{N} = 4$ SYM is a **gauged $\beta\gamma$ system**
[Beem, Lemos, Liendo, Peelaers, Rastelli, Rees '13]:

Symplectic bosons X, Y in the adjoint of $U(N)$:

$$X_b^a(z)Y_d^c(w) \sim \delta_d^a \delta_b^c \frac{1}{N} \frac{1}{z-w}$$

For the future, define a linear combination:

$$Z(u; z) \equiv X(z) + uY(z)$$

BRST charge:

$$Q_{\text{BRST}} \sim N \oint \text{Tr} \left(c[X, Y] + \frac{1}{2} b[c, c] \right)$$

B-model/Kodaira-Spencer theory

- The chiral algebra \mathcal{A}_N is dual to **topological B-model on $SL(2, \mathbb{C})$**
- The spacetime theory is **Kodaira-Spencer (BCOV) theory**
[Bershadsky, Cecotti, Ooguri, Vafa '93 '94]
 - holomorphic, depends only on complex structure (no metric)
- The chiral algebra \mathcal{A}_N is supported by N B-branes wrapping $\mathbb{C} \subset \mathbb{C}^3$

”**Derivation**” [Costello, Gaiotto '18]:

B-model on $\mathbb{C}^3 + N$ D1-branes \longrightarrow B-model on $SL(2, \mathbb{C}) \approx \text{AdS}_3 \times S^3$

\uparrow
 \mathcal{A}_N

Holographic dictionary

Chiral algebra \mathcal{A}_N	B-model on $SL(2, \mathbb{C})$
single traces	local modifications of asymptotic boundary condition
(sub)determinants	Giant Graviton D1-branes

In particular, determinant operator

$$\det(m + Z(u; z)), \quad Z(u; z) = X(z) + uY(z)$$

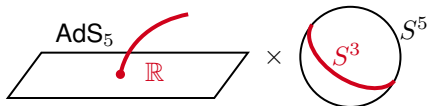
is dual to

- z = position at the boundary of AdS_3
- u = orientation of $S^1 \subset S^3$
- m controls size of $S^1 \subset S^3$



Giant Gravitons

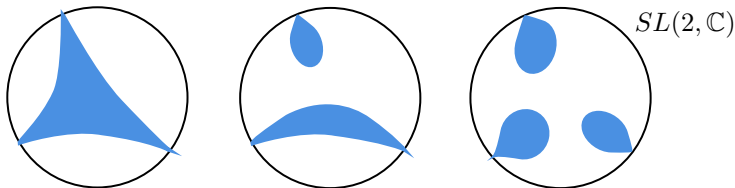
Determinant operators in $\mathcal{N} = 4$ SYM are dual to Giant Graviton D3-branes in $\text{AdS}_5 \times S^5$:



Future direction: "bootstrap" D3-branes from B-model D1-branes

Determinants and Giant Gravitons

Many possible brane configurations with the same boundary behaviour



We will match saddles ρ^* of correlation functions of determinants with brane configurations

- m_i, u_i, z_i control boundary behaviour
- saddles ρ^* will control the shape in the bulk

Determinant correlation functions

[Jiang, Komatsu, Vescovi '19]

see also [Chen, de Mello Koch, Kim, Van Zyl '19] [Berenstein, Wang '22]

- Fermionize determinants

$$\mathcal{D}(m; u; z) \equiv \det(m + Z(u; z)) = \int [d\bar{\psi}d\psi] e^{\bar{\psi}(m+Z(u,z))\psi}$$

- Rewrite correlation function using auxiliary bosonic variables ρ_j^i for $i \neq j$,
 $\rho_i^i \equiv m_i$

$$\left\langle \prod_i \mathcal{D}(m_i; u_i; z_i) \right\rangle \sim \int [d\rho] e^{\frac{N}{2} \sum_{i \neq j} \frac{z_i - z_j}{u_i - u_j} \rho_j^i \rho_i^j} (\det \rho)^N$$

- Saddle point analysis

$$\left\langle \prod_i \mathcal{D}(m_i; u_i; z_i) \right\rangle \sim \int [d\rho] e^{N S[\rho]}$$

with action

$$S[\rho] = \frac{1}{2} \sum_{i \neq j} \frac{z_i - z_j}{u_i - u_j} \rho_j^i \rho_i^j + \log \det \rho$$

Saddles and branes

Saddle point equations in the matrix form:

$$[\zeta, \rho] + [\mu, \rho^{-1}] = 0$$

where

$$\zeta = \begin{pmatrix} z_1 & & \\ & \ddots & \\ & & z_k \end{pmatrix}, \quad \mu = \begin{pmatrix} u_1 & & \\ & \ddots & \\ & & u_k \end{pmatrix}, \quad \rho = \begin{pmatrix} m_1 & & ? \\ & \ddots & \\ ? & & m_k \end{pmatrix}$$

We will match saddles ρ^* to classical brane configurations in B-model on $SL(2, \mathbb{C})$.

For each ρ^* we will define a **spectral curve** S_{ρ^*} in $SL(2, \mathbb{C})$ and check it matches dual Giant Graviton brane.

Spectral curve

For each saddle ρ^* we define a spectral curve S_{ρ^*} :

- Define **commuting matrices**:

$$B(a) = a\mu - \rho, \quad C(a) = a\zeta + \rho^{-1}, \quad D(a) = a\zeta\mu + \rho^{-1}\mu - \zeta\rho,$$

which satisfy

$$aD(a) - B(a)C(a) = 1.$$

- Define spectral curve:

$$\mathcal{S}_{\rho^*} = \{(a, b, c, d)$$

s.t. b, c, d are **simultaneous eigenvalues** of $B(a), C(a), D(a)\}$

- \mathcal{S}_{ρ^*} comes with a natural line bundle :

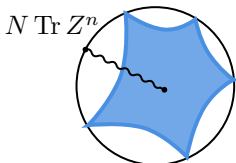
$$\mathcal{L}_{\rho^*} = \text{common eigenline of } B(a), C(a), D(a)$$

Holographic checks

Boundary behaviour $a \rightarrow \infty$:

$$\frac{B(a)}{a} = \begin{pmatrix} u_1 - \frac{m_1}{a} & & \\ & \ddots & \\ & & u_k - \frac{m_k}{a} \end{pmatrix} + \dots, \quad \frac{C(a)}{a} = \begin{pmatrix} z_1 + \frac{p_1}{a} & & \\ & \ddots & \\ & & z_k + \frac{p_k}{a} \end{pmatrix} + \dots,$$

where $p_i \equiv \partial S / \partial m_i = [\rho^{-1}]_i^i$.



Various holographic checks:

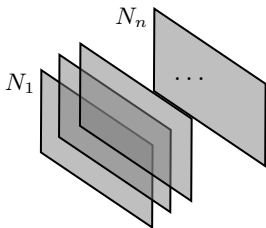
- Correlation functions of determinants with a **single-trace**
- Modifications of determinant / excitations of the brane

Coulomb branch geometries

- Duality can be extended to **non-conformal vacua** of the chiral algebra \mathcal{A}_N
- Twisted analog of [Kraus, Larsen, Trivedi '98]:

Coulomb branch of $\mathcal{N} = 4$ SYM \longleftrightarrow multi-center solutions

- Dual Calabi-Yau geometries are **deformations of $SL(2, \mathbb{C})$**
- Holographic check:
 - Determinant correlation functions (with a single-trace) and dual Giant Graviton branes



Vacua of chiral algebras

Translation-invariant vacuum \mathcal{V} of a chiral algebra: vevs $\langle \mathcal{O}_a(0) \rangle_{\mathcal{V}}$

$$\langle \mathcal{O}_a(z) \mathcal{O}_b(0) \rangle_{\mathcal{V}} = \sum_n z^{-n} \langle [\mathcal{O}_a \mathcal{O}_b]_n(0) \rangle_{\mathcal{V}}$$

By cluster decomposition for $z \rightarrow \infty$:

$$\langle [\mathcal{O}_a \mathcal{O}_b]_n(0) \rangle_{\mathcal{V}} = 0, \quad n \leq -1$$

The quotient by the above is **the C_2 -algebra of the VOA** [Zhu '96]

It is a Poisson algebra with product $[\mathcal{O}_a \mathcal{O}_b]_0$, which commutes with taking a vev:

$$\langle [\mathcal{O}_a \mathcal{O}_b]_0(0) \rangle_{\mathcal{V}} = \langle \mathcal{O}_a(0) \rangle_{\mathcal{V}} \langle \mathcal{O}_b(0) \rangle_{\mathcal{V}}$$

This gives a map:

space of vacua \mathcal{V} \longleftrightarrow **reduced spectrum of the C_2 -algebra**

Vacua of chiral algebras

- **Physical interpretation:**

The reduced spectrum of the Zhu's C_2 -algebra (the associated variety of a VOA) is the space of (translation-invariant) vacua of the VOA.

- In case of **chiral algebras of $\mathcal{N} = 2$ SCFT** [Beem, Rastelli '17]:

Higgs branch of $\mathcal{N} = 2$ SCFT \leftrightarrow the associated variety of the VOA

- Higgs branch of $\mathcal{N} = 2$ SCFT maps to the moduli space of its chiral algebra subsector

Future directions

- Spectral curve construction in other examples of twisted or free field holography
- Find SUSY D3-branes in $\text{AdS}_5 \times S^5$ that correspond to our B-model D1-branes
- Consider $\mathcal{O}(N^2)$ operators eg. $(\det Z)^N$, which are dual to backreacted geometries
- Analyze which saddles actually contribute
- Holographic dual of the **holomorphic twist of $\mathcal{N} = 4$ SYM**
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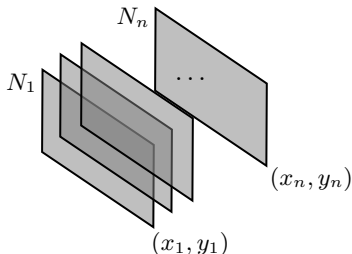
Thank you!

Coulomb branch geometries

Coordinates (x, y, z) in \mathbb{C}^3

N_i parallel B-branes wrapping \mathbb{C} 's at

$$x = x_i, \quad y = y_i \quad i = 1, \dots, n$$



Deformed geometry can be described in patches: $\forall i \ x \neq x_i$ or $y \neq y_i$

Coordinate transformation between patches I and I' :

$$z_I = z_{I'} + \frac{N_i/N}{(x - x_i)(y - y_i)}$$

For standard $SL(2, \mathbb{C})$ geometry:

$$z_0 - z_\infty = \frac{1}{xy}$$